

# STATISTICAL MECHANICS

## Final Exam – Spring 2003

**PROBLEM #1 :** (10 points) Use the grand canonical ensemble to find the expression of the entropy for a gas of non-interacting particles obeying (i) Bose-Einstein statistics, and (ii) Fermi-Dirac statistics in terms of the mean occupation number of a single-particle state with energy  $\epsilon$ ,  $\langle n_\epsilon \rangle$ .

*Hint:* In the grand canonical ensemble the entropy is given by

$$S = kT \left( \frac{\partial \ln(\mathcal{Z})}{\partial T} \right)_{z,V} - \langle N \rangle k \ln(z) + k \ln(\mathcal{Z})$$

where  $\mathcal{Z}$  and  $\langle N \rangle$  are the grand partition function and the expectation value of the total number of particles, respectively.

**PROBLEM #2 :** (12 points) Consider a two-dimensional ideal Bose gas. Let  $V = L^2$  be the area available to the system. The number of particles (which is conserved) is given by

$$N = z \frac{\partial}{\partial z} \ln[\mathcal{Z}(z, V, T)] = \sum_{\mathbf{p}} [z^{-1} \exp(\beta \epsilon_{\mathbf{p}}) - 1]^{-1} ,$$

where  $\mathcal{Z}$  is the grand partition function.

- Consider the system in the thermodynamic limit and discuss the importance of the  $\mathbf{p} = 0$  state.
- Show that there is no Bose-Einstein condensation at  $T \neq 0$  for a two-dimensional ideal gas.
- What is the groundstate of the system?

**PROBLEM #3 :** (8 points) The atom  $He^3$  consists of three nucleons, has spin 1/2 and is a fermion. The density of liquid  $He^3$  is  $0.081 \text{g/cm}^3$  near absolute zero. Calculate the Fermi energy  $\epsilon_F$  (in  $meV$ ) and the Fermi temperature  $T_F$  (in kelvin).

$\hbar = 1.055 \times 10^{-27} \text{erg/s}$ ,  $m_{nucleon} = 1.67 \times 10^{-24} \text{g}$ ,  $k_B = 1.38 \times 10^{-16} \text{ergK}^{-1}$ ,  $1eV = 1.6022 \times 10^{-12} \text{erg}$ .

# Final Exam - Solutions

Problem #1: Entropy

$$S = kT \left( \frac{\partial \ln \mathcal{Z}}{\partial T} \right)_{z, V} - \bar{N} k \ln z + k \ln \mathcal{Z}$$

$$\ln \mathcal{Z} = \frac{1}{a} \sum_{\epsilon} \ln(1 + a z e^{-\beta \epsilon})$$

$$S = k \sum_{\epsilon} \left\{ \frac{T}{a} \frac{\partial}{\partial T} \ln(1 + a z e^{-\beta \epsilon}) \Big|_{z, V} - \langle n_{\epsilon} \rangle \ln z + \frac{1}{a} \ln(1 + a z e^{-\beta \epsilon}) \right\}$$

$$= k \sum_{\epsilon} \left\{ \frac{\epsilon}{kT} \frac{1}{z e^{\beta \epsilon} + a} - \langle n_{\epsilon} \rangle \ln z + \frac{1}{a} \ln(1 + a z e^{-\beta \epsilon}) \right\}$$

$$\langle n_{\epsilon} \rangle = \frac{1}{z e^{\beta \epsilon} + a} \quad \rightarrow \quad \beta \epsilon - \ln z = \ln \frac{1 - a \langle n_{\epsilon} \rangle}{\langle n_{\epsilon} \rangle}$$

$$1 + a z e^{-\beta \epsilon} = \frac{1}{1 - a \langle n_{\epsilon} \rangle}$$

$$S = k \sum_{\epsilon} \left\{ \langle n_{\epsilon} \rangle \ln \frac{1 - a \langle n_{\epsilon} \rangle}{\langle n_{\epsilon} \rangle} + \left(-\frac{1}{a}\right) \ln(1 - a \langle n_{\epsilon} \rangle) \right\}$$

$$= k \sum_{\epsilon} \left\{ (\langle n_{\epsilon} \rangle - \frac{1}{a}) \ln(1 - a \langle n_{\epsilon} \rangle) - \langle n_{\epsilon} \rangle \ln \langle n_{\epsilon} \rangle \right\}$$

BE:  $a = -1$

$$S = k \sum_{\epsilon} \left\{ (\langle n_{\epsilon} \rangle + 1) \ln(1 + \langle n_{\epsilon} \rangle) - \langle n_{\epsilon} \rangle \ln \langle n_{\epsilon} \rangle \right\}$$

FD:  $a = +1$

$$S = k \sum_{\epsilon} \left\{ -(1 - \langle n_{\epsilon} \rangle) \ln(1 - \langle n_{\epsilon} \rangle) - \langle n_{\epsilon} \rangle \ln \langle n_{\epsilon} \rangle \right\}$$

## Problem # 2 :

$$(a) \quad N = \frac{L^2}{(2\pi\hbar)^2} \int_0^\infty dp \, 2\pi p \frac{1}{z^{-1} e^{\beta p^2/2m} - 1} + \frac{z}{1-z}$$

$$\frac{N}{L^2} = \frac{2\pi m\hbar T}{h^2} \int_0^\infty dx \frac{1}{z^{-1} e^x - 1} + \frac{1}{L^2} \frac{z}{1-z} = \frac{2\pi m\hbar T}{h^2} g_1(z) + \frac{1}{L^2} \frac{z}{1-z} = n$$

if  $z < 1$  the fraction of bosons in the  $\vec{p}=0$  state is zero in the thermodynamic limit

if  $z \rightarrow 1$  the fraction of bosons in the  $\vec{p}=0$  state is now finite.

(b) Evaluate  $g_1(z)$  :  $z < 1$

$$\int_0^\infty dx \frac{1}{z^{-1} e^x - 1} = - \int_0^\infty dx \frac{e^{-x}}{e^{-x} - z^{-1}} = - \int_0^1 dy \frac{1}{y - z^{-1}} = - \ln(1-z) = g_1(z)$$

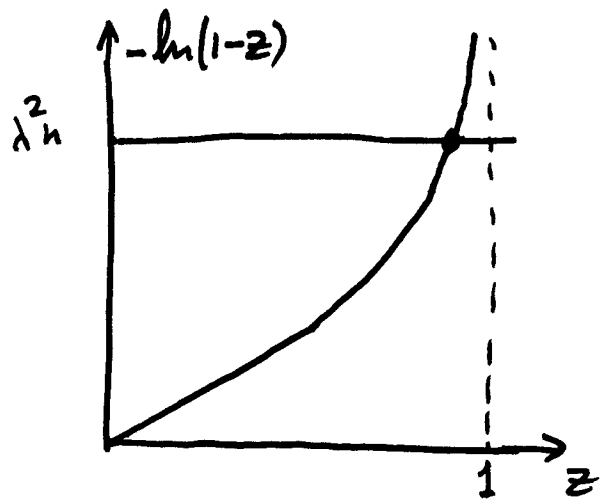
$$n = \frac{2\pi m\hbar T}{h^2} (-\ln(1-z)) + \frac{1}{L^2} \frac{z}{1-z} \quad ; \quad \lambda^2 = \frac{h^2}{2\pi m\hbar T} \text{ thermal length}$$

$$-\ln(1-z) = \lambda^2 n - \frac{\lambda^2}{L^2} \frac{z}{1-z}$$

In the thermodynamic limit ( $L \rightarrow \infty$ ) and for  $z < 1$  there is only one solution for given  $\lambda^2 n$ .

Hence, for  $T \neq 0$ ,  $\lambda^2 n$  is finite and  $z < 1$ . No condensation for

$T > 0$ .



(c) At  $T=0$  all particles are in the  $\vec{p}=0$  state. The ground state consists of all particles in the condensate. Hence,  $T_c = 0$ !

Problem #3: liquid  $\text{He}^3$

density:  $\rho = 0.081 \frac{\text{g}}{\text{cm}^3} = \frac{N}{V} m$

Fermi sphere:  $N = 2 \sum_{\mathbf{k}} 1_{\text{occ}} = \frac{2V}{(2\pi)^3} 4\pi \int_0^{k_F} dk k^2 = \frac{V}{(2\pi)^3} \frac{8\pi}{3} k_F^3 = V \frac{k_F^3}{3\pi^2}$

$$k_F = \left( \frac{3\pi^2 N}{V} \right)^{1/3} = (3\pi^2 n)^{1/3}$$

Fermi energy: 
$$\begin{aligned} E_F &= \frac{\hbar^2 k_F^2}{2m} = \frac{\hbar^2}{2m} (3\pi^2 n)^{2/3} \\ &= \frac{(1.055 \times 10^{-27} \text{ erg}\cdot\text{s})^2}{2 \times 3 \times 1.67 \times 10^{-24} \text{ g}} \cdot \left( 3\pi^2 \frac{\rho}{m} \right)^{2/3} \\ &= \frac{(1.055 \times 10^{-27} \text{ erg}\cdot\text{s})^2}{2 \times 3 \times 1.67 \times 10^{-24} \text{ g}} \left( 3\pi^2 \frac{0.081 \text{ g}/\text{cm}^3}{3 \times 1.67 \times 10^{-24} \text{ g}} \right)^{2/3} \\ &= 6.85 \times 10^{-16} \text{ erg} \approx 0.43 \text{ meV} \end{aligned}$$

Fermi temperature:  $T_h = \frac{E_F}{k_B} \approx 5 \text{ K}$