

QUANTUM MECHANICS

Homework set #20:

Relativistic Quantum Mechanics

Problem # 20.1 : One-dimensional Dirac equation with square-well potential

Consider the Hamiltonian for a Dirac particle of mass m moving in one spatial dimension under the influence of a scalar "square-well" potential,

$$H = -i\alpha \frac{d}{dx} + \beta[m + S(x)] ,$$

where $\alpha = \sigma_y$, $\beta = \sigma_z$, and the attractive scalar potential is given by

$$S(x) = \begin{cases} -S_0 & \text{for } |x| < L/2 \text{ ,} \\ 0 & \text{for } |x| > L/2 \text{ .} \end{cases}$$

Use units in which $\hbar = c = 1$ throughout this problem.

- (a) Show that if $\psi(x)$ is a solution of the Dirac equation ($H\psi = E\psi$) then so is the parity-transformed state, $\psi_p(x) \equiv \sigma_z\psi(-x)$. Recall that Pauli matrices anticommute.
- (b) Show that if $\psi(x)$ is a solution of the Dirac equation with energy $+E$, then the charge-conjugate state, $\psi_c(x) \equiv \sigma_x\psi(x)$, is also a solution, but with energy $-E$.
- (c) Propose a solution of the form,

$$\psi(x) = \begin{pmatrix} F(x) \\ G(x) \end{pmatrix} ,$$

and find the set of coupled first-order differential equations satisfied by F and G . Conclude that while F and G are continuous at $x = \pm L/2$, the first derivatives of both functions are not.

- (d) Assuming that there is no degeneracy, state the conditions on F and G so that a solution with definite parity, i.e. $\psi_p(x) = \pm\psi(x)$, might be found.
- (e) Using your results from (c) and (d), find the conditions for which a positive-parity, positive-energy bound-state ($0 \leq E < m$) solution exists. Note that you should state

these conditions in the form of transcendental equations, but do not attempt to solve the equations.

Problem # 20.2 : Dirac equation for a massless particle

Consider the Dirac equation in Hamiltonian form for a massless particle, e.g. using units in which $\hbar = c = 1$,

$$i \frac{\partial}{\partial t} \psi(\vec{x}, t) = (-i \vec{\alpha} \cdot \vec{\nabla}) \psi(\vec{x}, t) ,$$

where $\vec{\alpha} = \alpha_x \hat{x} + \alpha_y \hat{y} + \alpha_z \hat{z}$, and the α_i are matrices.

- (a) Find the anticommutation relations that α_i must satisfy if each component of $\psi(\vec{x}, t)$ is required to satisfy the Klein-Gordon equation,

$$\left(\frac{\partial^2}{\partial t^2} - \vec{\nabla}^2 \right) \psi(\vec{x}, t) = 0 .$$

- (b) Show that the 2×2 matrices $\alpha_i = \pm \sigma_i$ satisfy the conditions found in part (a). The σ_i are Pauli matrices,

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} , \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} , \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} .$$

- (c) Look for satisfactory solutions to the Dirac equation when $\alpha_i = +\sigma_i$, of the form

$$\psi_+(\vec{x}, t) = \begin{pmatrix} F_+ \\ G_+ \end{pmatrix} \exp(-iEt + i\vec{p} \cdot \hat{x}) ,$$

where \vec{p} is the momentum of the particle and $E = |\vec{p}|$. Find the equations satisfied by $\begin{pmatrix} F_+ \\ G_+ \end{pmatrix}$.

- (d) Look for stationary solutions to the Dirac equation when $\alpha_i = -\sigma_i$, of the form

$$\psi_-(\vec{x}, t) = \begin{pmatrix} F_- \\ G_- \end{pmatrix} \exp(-iEt + i\vec{p} \cdot \hat{x}) ,$$

where \vec{p} is the momentum of the particle and $E = |\vec{p}|$. Find the equations satisfied by $\begin{pmatrix} F_- \\ G_- \end{pmatrix}$.

- (e) Explain which of the solutions of parts (c) and (d) would be appropriate to describe neutrinos, which are massless particles with negative helicity.

Problem # 20.3 : Zitterbewegung

Consider a free Dirac particle described by

$$H = c\vec{\alpha} \cdot \vec{p} + \beta mc^2 ,$$

where $\vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}$ and $\beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$. The matrices $\vec{\sigma}$ and 1 are the 2×2 Pauli and unit matrices.

- (a) Use the Heisenberg picture to establish

$$\frac{d}{dt}\vec{p} = 0 \quad , \quad \frac{d}{dt}\vec{\alpha} = \frac{2}{i\hbar}(c\vec{p} - H\vec{\alpha}) \quad , \quad \frac{d}{dt}\vec{r} = c\vec{\alpha} .$$

- (b) These equations can be integrated to obtain

$$\begin{aligned} \vec{p} &= \vec{p}(0) & \vec{\alpha} &= \frac{c\vec{p}}{H} + e^{2iHt/\hbar} \left[\vec{\alpha}(0) - \frac{c\vec{p}}{H} \right] \\ \vec{r} &= \vec{r}(0) + \frac{c^2\vec{p}}{H}t + \frac{\hbar c}{2iH} \left\{ e^{2iHt/\hbar} - 1 \right\} \left[\vec{\alpha}(0) - \frac{c\vec{p}}{H} \right] . \end{aligned}$$

Confirm this by differentiating the given results for \vec{p} , $\vec{\alpha}$, and \vec{r} .

- (c) The last terms in the result for $\vec{\alpha}$ and \vec{r} are related to the so-called "Zitterbewegung". These make no contribution to $\langle \vec{\alpha} \rangle$ and $\langle \vec{r} \rangle$ for wave packets containing solely positive or negative energy components. Prove this by establishing

$$\Lambda_{\pm} \left(\frac{\vec{p}}{H} - \vec{\alpha} \right) \Lambda_{\pm} = 0 ,$$

where Λ_{\pm} are the projectors on states with positive and negative energy, respectively,

$$\Lambda_{\pm} = \frac{1}{2} \left[1 \pm \frac{H}{E_p} \right] ,$$

with $E_p = \sqrt{m^2c^4 + p^2c^2}$.

Hint: Start by establishing the general result $\Lambda[\Lambda, O]\Lambda = 0$ for any projection Λ and operator O , then consider $O = \vec{\alpha}$.

Problem # 20.4 : Spin-orbit coupling

The spin-orbit interaction in the hydrogen atom is given by the potential

$$V_{so} = \frac{1}{2m^2c^2} \frac{1}{r} \frac{dV}{dr} \vec{S} \cdot \vec{L} ,$$

which is of relativistic origin (i.e. it vanishes as $c \rightarrow \infty$). Here \vec{S} and \vec{L} are the spin and orbital angular momenta, respectively, and $V = -Ze^2/r$ is the Coulomb potential.

Using the wave-functions for the hydrogen atom as given by the solutions of the Schrödinger equation (non-relativistic hydrogen atom), obtain the spin-orbit coupling constant by calculating the matrix element

$$\langle nlm_l m_s | V_{so} | nlm'_l m'_s \rangle .$$

Obtain the energies of the hydrogen atom including the spin-orbit coupling and compare your result to the exact energies (solutions of the Dirac eqn.), given by

$$E_{nj} = -\frac{mc^2}{2} \frac{\alpha^2 Z^2}{n^2} \left[1 + \frac{\alpha^2 Z^2}{n^2} \left(\frac{n}{j + 1/2} - \frac{3}{4} \right) + \dots \right] .$$

Here $\alpha = e^2/\hbar c \approx 1/137$ is the fine-structure constant.

Hint: You will need the expectation value

$$\langle nl | r^{-3} | nl \rangle = \frac{Z^3}{a_0^3 n^3} \frac{1}{l(l + 1/2)(l + 1)} , \quad l \neq 0 ,$$

which can be calculated using the method of problem # 8.1 (calculation not required).

Here $a_0 = \hbar^2/me^2$ is the Bohr radius for $Z = 1$.