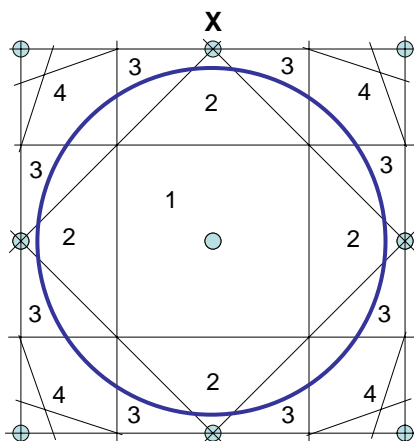


Condensed Matter Physics I — PHZ 5491

Midterm II November 15, 2007

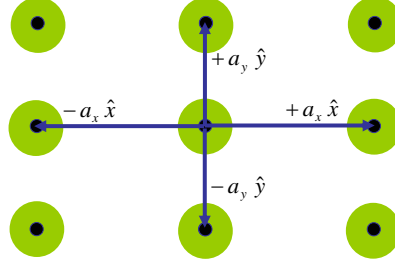
1. (50 pts) Consider a two-dimensional metal consisting of a gas of electrons moving in a weak periodic potential due to a square lattice with lattice constant a . Assume the number density of electrons is $n = 5/a^2$, i.e. five electrons per primitive cell.

- (a) Determine the Fermi wave vector k_F for free electrons in two dimensions with number density $n = 5/a^2$.
- (b) From the result of (a) verify that the free electron Fermi surface is entirely contained within the first four Brillouin zones. To do this it is sufficient to show that k_F does not quite reach the point labeled **X** in the figure below. (In this figure the blue dots are reciprocal lattice points and numbers label the different Brillouin zones.)



- (c) In the reduced zone scheme carefully sketch the nearly free electron Fermi surfaces in the first four zones. In each case indicate which states are occupied and which states are unoccupied.
- (d) If a magnetic field is applied in the z direction (i.e. perpendicular to the two-dimensional gas) indicate for each of the zones in Part (c) whether the carriers contribute positively or negatively (or zero in the case of a filled band) to n_{eff} , the effective carrier density which determines the Hall coefficient. Do you expect the total n_{eff} to be positive or negative?

2. (50 pts) Consider a two-dimensional monatomic solid in which the atoms form a rectangular Bravais lattice with primitive vectors $\vec{a}_1 = a_x \hat{x}$ and $\vec{a}_2 = a_y \hat{y}$ where $a_x > a_y$ (see figure).



The tight-binding s -band for this lattice is determined by the matrix elements

$$\gamma(\vec{R}) = - \int d\mathbf{r} \phi_s(\vec{r}) \Delta U(\vec{r}) \phi_s(\vec{r} - \vec{R}),$$

where $\phi_s(\vec{r})$ is an atomic s -orbital associated with the atom at the origin and $\Delta U(\vec{r})$ is the perturbation due to the other ions in the lattice.

By symmetry we know that $\gamma(+a_x \hat{x}) = \gamma(-a_x \hat{x})$ and $\gamma(+a_y \hat{y}) = \gamma(-a_y \hat{y})$ and so we can define γ_x and γ_y as follows:

$$\gamma_x = \gamma(\pm a_x \hat{x}); \quad \gamma_y = \gamma(\pm a_y \hat{y}).$$

(a) Show that the tight-binding s -band for this system has the form

$$\mathcal{E}(\vec{k}) = \mathcal{E}_0 - 2\gamma_x \cos k_x a_x - 2\gamma_y \cos k_y a_y.$$

(b) Taylor expand $\mathcal{E}(\vec{k})$ about the point $\vec{k} = 0$ and show that near this point the dispersion is quadratic and can be expressed in the following form,

$$\mathcal{E}(\vec{k}) \simeq \text{Const.} + \frac{\hbar^2}{2} (k_x, k_y) \cdot \mathbf{M}^{-1} \cdot \begin{pmatrix} k_x \\ k_y \end{pmatrix}.$$

Determine the matrix elements of the (2×2) inverse effective mass tensor \mathbf{M}^{-1} in terms of γ_x , γ_y , a_x , a_y and \hbar .

(c) Now imagine a magnetic field \vec{B} is applied in the positive z direction (out of the page). Sketch a typical cyclotron orbit in k space near the point $\vec{k} = 0$ where the quadratic dispersion given above is valid. On this sketch clearly indicate the direction in which the orbit is traversed.

(d) Obtain the period of this cyclotron orbit as a function of γ_x , γ_y , a_x , a_y , \hbar , c , e and B .